Ultrasound contrast agents modeling using an extended Volterra model

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Harmonic imaging has historically occurred with the introduction of ultrasound contrast agents such as microbubbles. These agents, due to their nonlinear behavior, have provided a great increase in the contrast of ultrasound images. Although this modality has revolutionized the clinical practice, it still suffers from the presence of harmonics in the echo of the tissue that reduces its efficiency. One way to overcome this problem was to turn to the sub- and ultra-harmonic imaging based on the reception of sub-harmonics \((\frac{1}{2}f_0)\) and ultra-harmonics \((\frac{2}{3}f_0, \frac{5}{3}f_0,...)\) generated by the microbubbles at high pressure levels. Nonlinear existing models like Volterra series are able to model harmonics only. In this paper, we proposed an extended Volterra model able to model and extract these sub- and ultra-harmonics. Results showed that sub- and ultra-harmonics were extracted and separated from other harmonic components. Signals backscattered by the medium infused with microbubbles can be accurately represented by the derived model. The gain achieved with our method was 3.7 dB compared to the standard Volterra modeling. In the frequency domain, the spectrum of the simulated signal has described perfectly that of the microbubble.

1 Introduction

The diagnostic capabilities of ultrasound imaging have been greatly improved by the introduction of ultrasound contrast agents such as microbubbles [1]. This contrast enhancement is due to two principal reasons. From one hand side, the acoustic impedance of microbubbles is largely different than that of tissues [2, 3] and, from the other side, the response of microbubbles is nonlinear [2, 4]. This nonlinearity is directly observable through the presence of harmonics that can be extracted by filtering [5]. For instance, in second harmonic imaging, the principle is to excite the medium at a frequency \(f_0\) and to reconstruct the image at the first harmonic \(i.e.\) around \(2f_0\) [5, 6]. However, the generation of nonlinear components during the propagation in tissue limits the contrast. These nonlinear components can be reduced in part by reducing the emitting pressure [5, 6, 7, 8]. One alternative solution to distinguish microbubbles from surrounding tissue is to use the difference in signatures between the two media through sub- and ultra-harmonic imaging. Unlike tissues, experimental studies have shown the existence of sub- and ultra-harmonics in the response of microbubbles [1, 6, 8, 9]. In ultrasound contrast imaging, subharmonic imaging consists of emitting at a frequency \(f_0\) and receiving at the sub-harmonic \(\frac{1}{2}f_0\) while ultraharmonic imaging is based on the reception of \((\frac{2}{3}f_0, \frac{5}{3}f_0,...)\). Many in vitro studies showed the possibility of the application of subharmonic imaging in early detection of the angiogenesis around tumors [3]. The reception of sub-harmonics may be done using a narrow band transducer for the reception with center frequency \(\frac{1}{2}f_0\)[3, 10].

In order to improve the image contrast, it is necessary to understand the behavior of microbubbles. Indeed, the dynamics of the microbubble is quite well understood and many nonlinear models have been proposed [2]. Various modifications were performed to accommodate the microbubble shell and the nonlinearities, including the Rayleigh-Plesset modified equation [11, 12, 13].

Although these models accurately simulate the oscillation of microbubbles, other models that consider the system to be studied as a black box, have been used to optimally model microbubbles signals. The best known one is the Volterra series which has a capacity to describe the input-output behavior of nonlinear dynamic systems [14]. Volterra series have already been employed in ultrasound contrast imaging and contributed to increase the contrast of images [15]. These series, that are able to model and extract the harmonic components only [14], are power series of integer orders. However, there is no simple mathematical model able to extract sub- and ultra-harmonics from microbubbles signals. However the modeling of sub-harmonics using Multiple Input Single Output (MISO) Volterra series is discussed in [16]. This approach is unable to extract these components and separate them from other components.

We proposed in this study an original solution based on the Volterra series that enabled to model and extract sub- and ultra-harmonic components. The originality in this work resided in the simplicity of the proposed method that used existing mathematical operators.

2 Materials and Method

The purpose of our dissertation was to develop a Volterra series based model to model microbubbles signals in presence of sub- and ultra-harmonics of order \(\frac{1}{2}\) \((f_0 = \frac{2}{3}f_0,\) where \(n\) is an odd integer) and to extract them. The nonlinear system to be identified was the microbubble excited by a pressure wave \(x(n)\) (the input signal). By solving numerically the nonlinear differential equation of a microbubble, subjected to a pressure signal, the backscattered signal by the microbubbles \(y(n)\) (the output signal) was obtained (see Figure 1).

The dynamics of a microbubble were simulated by solving the Rayleigh-Plesset modified equation using Hoff’s method [11]. The incident wave to the microbubble was a sinusoidal signal apodized with Hanning window, of frequency \(f_0 = 4\) MHz, 1.6 MPa pressure, and 32 cycles. Under the previous frequency and pressure conditions, the oscillation of the microbubble is nonlinear including sub- and ultra-harmonics [8]. The backscattered signal was sampled \(f_s = 36\) MHz. The parameters of microbubbles are shown in table 1.

| Table 1: Parameters of microbubbles [11]. |
|--------------|-----------------|
| resting radius | \(r_0 = 2\mu m\) |
| shell thickness | \(d_{Se} = 4nm\) |
| shear modulus | \(G_s = 50\) MPa |
| shear viscosity | \(\eta = 0.8\) Pa.s |

2.1 Volterra model

The Volterra model is equivalent to a power series decomposition of integer orders. To model and extract the sub- and ultra-harmonic frequencies, we proposed to extend the standard Volterra series by adding a component dedicated to them.
Volterra series were introduced as Taylor series with memory. Let \( y(n) \) be the scattered signal by the medium with microbubbles. The output \( y(n) \) can be expressed in terms of the present and previous inputs \( x(n) \) and \( x(n - k) \), respectively. The general form of the Volterra model \((3, m)\), of order 3 and memory \( m \), in the discrete time was as follows [17]:

\[
\hat{y}(n) = h_0 + \sum_{k=1}^{m-1} h_1(k_1) x(n - k_1)
+ \sum_{k_1=k_2=0}^{m-1} h_2(k_1, k_2) x(n - k_1) x(n - k_2)
+ \sum_{k_1=k_2=k_3=0}^{m-1} h_3(k_1, k_2, k_3) x(n - k_1) x(n - k_2) x(n - k_3)
\]

where \( \hat{y}(n) \) was the observed output sequence associated with the input sequence \( x(n) \), and \( h_p(k_1, k_2, \ldots, k_p) \) were the Volterra kernels of order \( p \) of the development. The order 3 of the Volterra series is sufficient given the limited bandwidth of ultrasound transducers [14]. If the function \( V^m_H \) were the Volterra model, \( \hat{y}_H(n) \) could be expressed as:

\[
\hat{y}_H(n) = V^m_H(x(n)).
\]

The Volterra’s coefficients \( h_p(k_1, k_2, \ldots, k_p) \) were obtained by minimizing the mean square error between the two signals \( y(n) \) and \( \hat{y}_H(n) \) (Figure 1):

\[
\text{arg min}[(y(n) - \hat{y}_H(n))^2].
\]

The calculation method was similar to that developed by [14] but without the Cholesky decomposition.

![Figure 1: Identification of the nonlinear system of the microbubble with the Volterra model.](image)

### 2.2 Extended Volterra model

The extended Volterra model that we proposed, allowed identifying the harmonic, sub-harmonic, and ultra-harmonic components. The corresponding output sequence of our extended Volterra model could be written as follows:

\[
\hat{y}(n) = V^m_{SUH}(x(n)).
\]

The solution that we proposed consisted of two parts:

\[
V^m_{SUH}(x(n)) = V^m_H(x(n)) + V^m_{SU}(x(n)).
\]

A part that modeled exclusively the harmonics:

\[
\hat{y}_H(n) = V^m_H(x(n)),
\]

and another that modeled the sub- and ultra-harmonics:

\[
\hat{y}_{SU}(n) = V^m_{SU}(x(n)).
\]

The reconstructed signal of the microbubble was then:

\[
\hat{y}(n) = \hat{y}_H(n) + \hat{y}_{SU}(n).
\]

### 2.3 Numerical procedure

The numerical procedure was written with Matlab (Mathworks, Natick, MA, USA). It included the following steps:

1. **Modeling of integer harmonics**: it was the Volterra model \((3, m)\). The obtained signal was given by equation (6);

2. **Modeling of sub- and ultra-harmonics of order \( \frac{1}{2} \)**:

   a) The analytic signal \( y_a(n) \), whose spectrum contained only positive frequencies was modulated by multiplying it by an exponential having a frequency \( \frac{1}{2} f_0 \). The modulated signal was then:

   \[
   y_{a_{mod}}(n) = y_a(n)e^{j2\pi \frac{f_0}{2} n},
   \]

   with:

   \[
   y_a(n) = y(n) + jH(y(n)),
   \]

   where \( H(y(n)) \) was the Hilbert transform of \( y(n) \). From a spectral point of view, the modulation shifted the spectrum by \( \frac{1}{2} f_0 \). Indeed, if the spectrum of \( y(n) \) included \( \frac{1}{2} f_0 \) and \( f_0 \) components, then the spectrum of \( y_{a_{mod}} \) will be composed of \( f_0 = (\frac{1}{2} f_0 + \frac{1}{2} f_0) \) and \( \frac{1}{2} f_0 = (\frac{1}{2} f_0 + \frac{1}{2} f_0) \) components. The signal \( y_{a_{mod}} \) contained shifted sub- and ultra-harmonics instead of harmonics.

   b) A Volterra model of order \((3, m)\) identified the real part of the modulated signal \( R(y_{a_{mod}}(n)) \) by minimizing the following relation:

   \[
   \text{arg min}[(R(y_{a_{mod}}(n)) - \hat{y}_2(n))^2].
   \]

   The Volterra model extracted the harmonic components around \( n f_0 \) (an integer) which were initially the sub- and ultra-components. The modeled signal became:

   \[
   \hat{y}_2(n) = V^m_{SUH}(x(n)).
   \]

   c) The analytic signal \( \hat{y}_2(n) \) was demodulated with a frequency \((-\frac{1}{2} f_0)\), that shifted the harmonic components as well as the sub- and ultra-harmonic components back to their original positions. The demodulated signal became:

   \[
   \hat{y}_{2_{demod}}(n) = \hat{y}_2(n)e^{-j2\pi \frac{f_0}{2} n}.
   \]

   The double application of the Hilbert transform introduced a negative sign. For this reason, we multiplied the demodulated signal by \( e^{j\pi} \). In order to obtain the positive and negative frequencies of the demodulated signal we considered its real part merely. We obtained:

   \[
   \hat{y}_{SU}(n) = R(\hat{y}_{2_{demod}}(n))e^{j\pi} = -R(\hat{y}_{2_{demod}}(n)).
   \]

The final modeled signal \( \hat{y}(n) \) was obtained by adding the two signals \( \hat{y}_H(n) \) and \( \hat{y}_{SU}(n) \) as given by equation (8). The extended Volterra model was symbolized by \((3, \frac{1}{2}, m)\) where 3 is the model order, \( \frac{1}{2} \) is the coefficient of modulation and \( m \) is the model memory.
of the extended Volterra model while figure 3 showed both systems $S_1$ and $S_2$.

The relative mean square error $RMSE$ used to evaluate the accuracy of our method was given by:

$$RMSE = \frac{|\hat{y}(n) - y(n)|^2}{|y(n)|^2}, \quad (15)$$

where $\hat{y}(n)$ was the final signal reconstructed by the extended Volterra model and $y(n)$ was the signal backscattered by the microbubble.

3 Results

To verify the efficiency of our method, we applied both the Volterra model $(3, 1, m)$ as well as the extended Volterra model $(3, \frac{1}{2}, m)$ to a synthetic microbubble signal. The microbubble was insonified with the sine pulse described in section 2 and the backscattered signal was then calculated.

Qualitatively, Figure 4(a) shows the backscattered signal $y(n)$ and that modeled by the extended Volterra model $(3, \frac{1}{2}, 4)$ $\hat{y}(n)$ versus time. Figure 4(b) shows the difference between $y(n)$ and $\hat{y}(n)$. Figure 4(c) shows the harmonic signal $\hat{y}_H(n)$ obtained with the standard model $(3, 4)$, and figure 4(d) shows the sub- and ultra-harmonic signal $\hat{y}_{SU}(n)$ extracted with the extended model $(3, \frac{1}{2}, 4)$.

Figure 5 shows the spectra of the various signals shown in figure 4. The spectral composition of these signals was in good qualitative agreement with figure 4. These figures show that sub- and ultra-harmonics components were modeled, extracted and separated from other harmonic components.

Quantitatively, table 2 shows a comparison of the $RMSE$ between the the signal modeled with the extended model and that modeled with the standard model. The gain achieved was $3.7$ dB (the difference between $RMSE$ between the reconstructed signal with the extended Volterra model and the backscattered signal $(-11.5)$ dB and the $RMSE$ between the reconstructed signal with the standard Volterra model and the backscattered signal $(-7.8)$ dB with the application of the extended Volterra model).
Table 2: Relative mean square errors $RMSE$ (dB) between the signal backscattered by the microbubble and that modeled with the Volterra model (3, 4), and between the signal backscattered by the microbubble and that modeled with the extended Volterra model (3, 4, 5).

<table>
<thead>
<tr>
<th>Model</th>
<th>Standard Volterra</th>
<th>extended Volterra</th>
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<tr>
<td>$RMSE$ (dB)</td>
<td>-7.8</td>
<td>-11.5</td>
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4 Discussion

The ultrasound contrast agents promoted the generation of sub- and ultra-harmonics that did not show up in the echo of tissues. The application of the extended Volterra model on simulated signals of microbubbles showed a high efficiency to extract and recover sub- and ultra-harmonics of order $\frac{1}{2} : f_n = \frac{1}{2} f_0$ apart of harmonics. This result overcame the weakness of the standard Volterra model which modeled harmonics only. Extracted sub- and ultra-harmonics could be used to make sub- and ultraharmonic contrast imaging in order to produce high contrast images.

In addition, the extended model modeled the microbubble signal with a gain 3.7 dB compared to the standard Volterra model. In the frequency domain, all the frequency components of the backscattered signal were correctly reconstructed. The identification process of the parameters enabled modeling of signals even at very high excitation pressure (1.6 MPa) known to exhibit sub-harmonics.

5 Conclusion

In our dissertation, we proposed an extended Volterra model. The originality of this model lied in enabling the extraction of sub- and ultra-harmonics components of microbubbles signals. In addition, modeling the microbubble signal, may be performed even at high pressure levels, unlike the standard Volterra model.

This method enabled the separation of sub- and ultra-harmonics components. Such tool could be greatly effective in the field of ultrasound contrast imaging. Especially the sub- and ultra-harmonic contrast imaging.

This work could be able to complete by calculating the parameters of sub- and ultra-harmonics in order to separate them, and developing an automatic method to calculate the contrast-to-tissue ratio (CTR) using these parameters.

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References